

Effect of frequency of air flow on the performance of the Wells turbine

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Experimental investigations on the effect of Strouhal number and rotor solidity on the performance of 0.2 m diameter Wells self rectifying air turbine with NACA 0021 profile blades are presented. The results show an increase in starting torque with the Strouhal number of air flow and rotor solidity. The effect of Strouhal number on the running performance of the turbine is solidity dependent

Keywords: turbomachinery, fluid flow, wave energy

The Wells air turbine is a self-rectifying air turbine suited for wave energy conversion from the oscillating water-air column devices. Some recent reports¹⁻³ explain the principle of the working of turbine, the starting characteristics, the prediction of the performance and the effect of several variables on the running performance of the Wells turbine. Some initial experiments on the Wells turbine in an oscillating air flow have been reported⁴. The starting and running performance of a 0.2 m Wells turbine at several air flow frequencies and two solidities are reported here. The results are correlated with the results from unidirectional (zero frequency) air flow tests.

Experiments

Experiments were performed in an oscillating flow rig (Fig 1) consisting of an air bellows driven by a ram and hydraulic jack. The frequency of the jack which produced air oscillations was controlled by a signal generator and servo-amplifier system, the vertical displacement of the piston being controlled by limit switches. The test section with a 0.2 m diameter Wells turbine was mounted on top of the test rig. The model configurations chosen for the tests were (i) $h=0.65$ and $S=0.69$ and (ii) $h=0.65$ and $S=0.52$. Both models had 0.09 m chord blades with NACA 0021 aerofoil sections. All the tests were carried out at $R=2.5 \times 10^5$ for running conditions and $R_x=5.5 \times 10^4$ for starting conditions.

The measurements consisted of velocity profiles on both sides of the rotor, pressure drops across the rotor, turbine torque and speed. Velocity profile traversing stations 1 and 2 are shown in Fig 1. A CTA anemometer (DISA type 55M10) with a lineariser (type 55D10) and a calibrated hot wire probe (P11) was used for the velocity measurements. The pressure drop was measured by a low differential pressure transducer (Furness Control type FC040) which had frequency response up to 50 Hz. The torque was measured by a torque transducer (ASL DORT1). An optical type tachometer was used for

the speed measurements. The data acquisition system consisted of a digital PAC 11/16 utilising a DEC LSI 11/2 microprocessor with a VT 50 graphic terminal and a Data Dynamics printer.

Tests were performed at several oscillating flow frequencies and several pressure drops. Some measurements were also made with the turbine rotor in a locked position (zero speed).

Running performance

Typical rotor pressure drop-time, ($\Delta p \sim t$) and inlet axial velocity-time ($V_x \sim t$) histories obtained during the tests were shown in Fig 2. These were close to sine waves. The variation of turbine torque τ and speed ω , (Fig 3) with time over the mean values was within $\pm 4\%$ and $\pm 2\%$ respectively. The power output variation with time was $\pm 6\%$. Fig 4 shows that at any given instant the inlet axial velocity was almost constant with the radius within the annulus. These figures are discussed in more detail elsewhere⁴. From the oscillating air flow tests the values of p_{RMS}^* , W_{av}^* , η_{av} and ϕ_{RMS} for a cyclic period T can be obtained as follows:

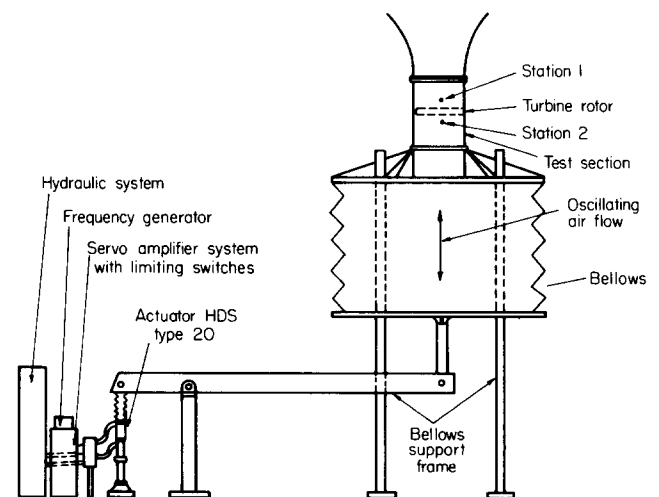


Fig 1 Oscillating air flow rig

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$$p_{RMS}^* = \frac{1}{\rho \omega_{av}^2 D_t^2} \left[\int_0^1 \Delta p^2 d\left(\frac{t}{T}\right) \right]^{1/2} \quad (1)$$

is the root mean square pressure drop;

$$W_{av}^* = \frac{1}{\rho \omega_{av}^3 D_t^5} \int_0^1 \omega \tau d\left(\frac{t}{T}\right) \quad (2)$$

is the nondimensional average power output;

$$\eta_{av} = \frac{\int_0^1 \omega \tau d\left(\frac{t}{T}\right)}{\int_0^1 \Delta p Q d\left(\frac{t}{T}\right)} \quad (3)$$

is the average cyclic efficiency, and:

$$\phi_{RMS} = \frac{2}{\omega_{av} D_t} \sqrt{\int_0^1 V_x^2 d\left(\frac{t}{T}\right)} \quad (4)$$

is the root mean square flow coefficient where:

$$\omega_{av} = \int_0^1 \omega d\left(\frac{t}{T}\right)$$

is the average speed, τ is the torque, and $R = U_t c / \nu$ is the Reynolds number based on the tip speed which, as the variation in speed was within $\pm 2\%$, was assumed constant.

Fig 5 shows plots of p^* versus ϕ_{RMS} for a range of oscillating flow frequencies from 0.2 to 1 Hz and for two solidities, 0.52 and 0.69. Also shown are the results from unidirectional flow tests (zero frequency). It can be observed that generally a linear relationship exists between p^* and ϕ_{RMS} . The effect of increasing airflow frequency on p^* is not very evident from this graph

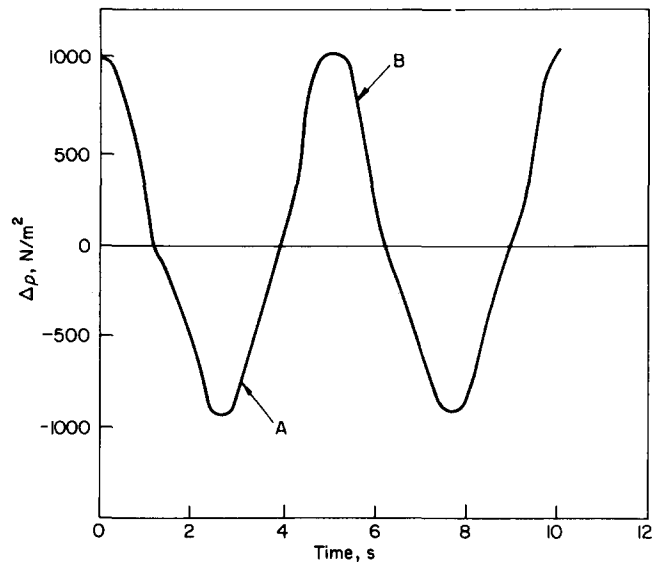
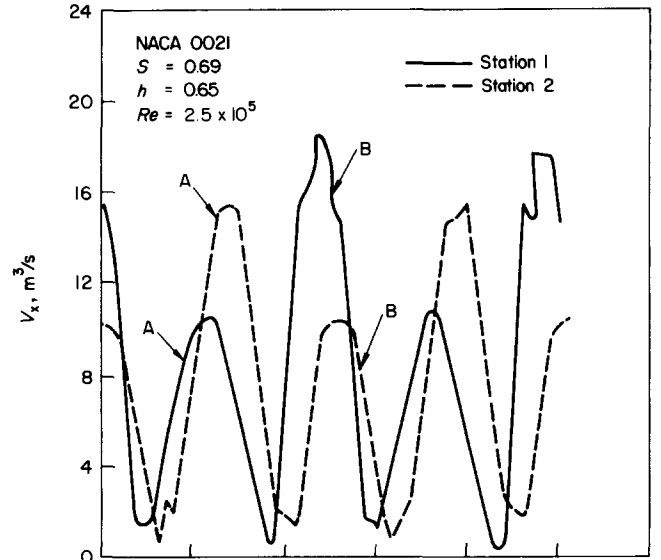


Fig 2 Pressure drop and axial velocity histories

Notation			
c	Blade chord length, m	W^*	Power coefficient, $W_0 / \rho \omega_{av}^3 D_t^5$
$C_{L,D}$	Lift, drag coefficient	U_t	Tip tangential velocity at rotor, m/s
$D_{t,h}$	Tip, hub diameter of the rotor, m	α	Angle of incidence = $\tan^{-1}(V_x / U_t)$, degrees
f	Frequency of oscillations, Hz	ρ	Density of air, kg/m ³
h	Hub to tip ratio D_h / D_t	ω	Angular velocity, rad/s
n	Number of blades	ν	Kinematic viscosity of air, m ² /s
$p^*, \Delta p$	Pressure drop coefficient $\frac{\Delta p}{\rho \omega_{av}^2 D_t^2}, \frac{\Delta p}{\frac{1}{2} \rho V_x^2}$	η	Gross efficiency W_0 / W_1
Δp	Pressure drop across rotor, N/m ²	$\tau, \bar{\tau}$	Torque Nm, reduced torque, $\tau / \frac{1}{2} \rho V_x^2 n c (D_t^2 / 4) (1 - h^2)$
Re, Re_x	Reynolds number $U_t c / \nu, V_x c / \nu$	τ^*	Torque coefficient, $\bar{\tau} / \Delta p$
S	Solidity $2nc / (\pi D_t (1 + h))$	ϕ	Flow coefficient (V_x / U_t)
S_t, S_x	Strouhal number $fc / U_t, fc / V_x$		
t	Time, s		
T	Period for a cycle,		
V_x	Axial velocity, m/s		
		Subscripts	
		av	Average value
		opt	Optimum value
		rms	Root mean square value

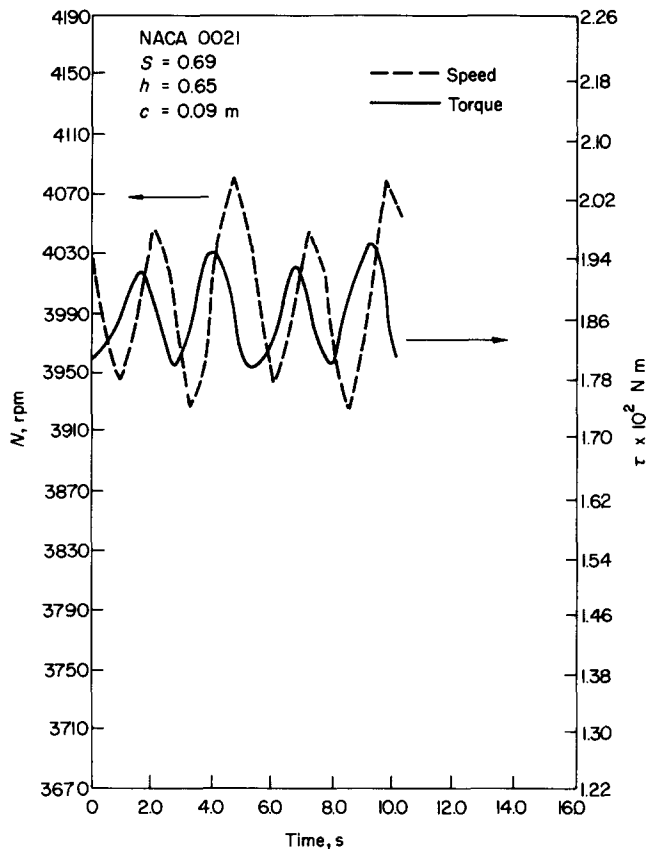


Fig 3 Torque and speed histories

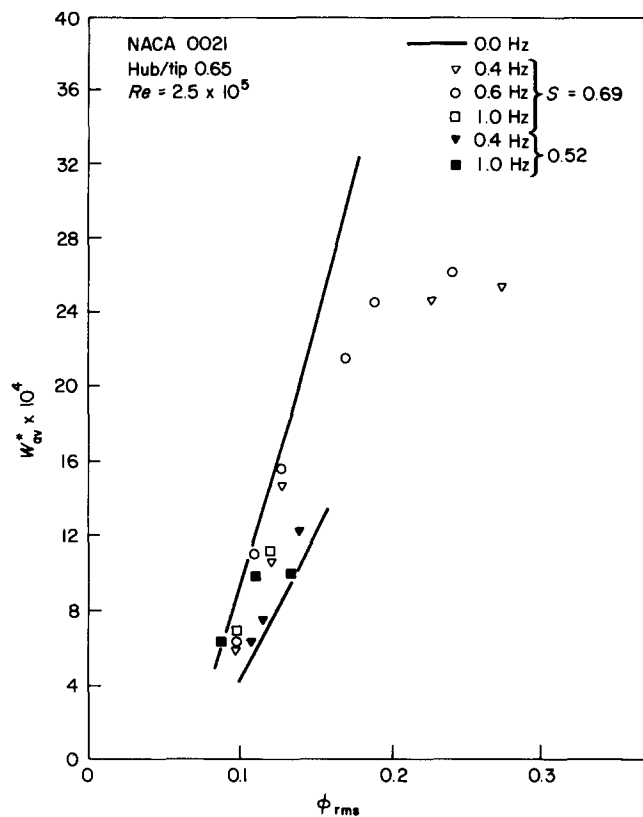


Fig 5 Effect of solidity on pressure drop

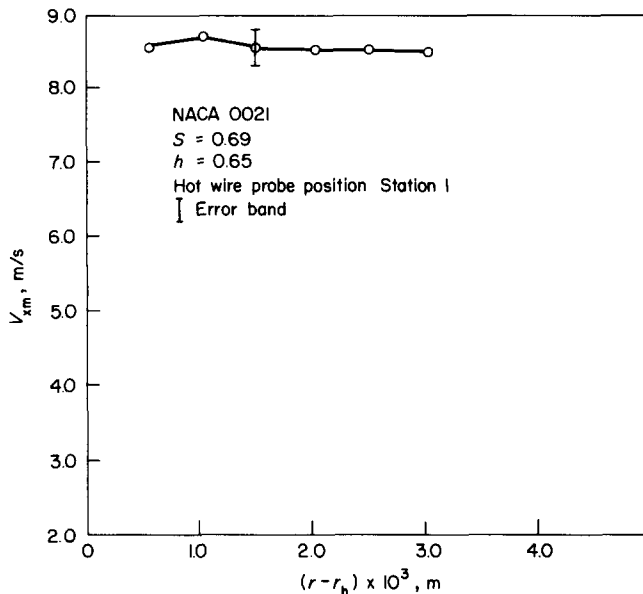


Fig 4 Annulus peak velocity distribution

although for a solidity of 0.52 the oscillating flow test results show a reduction in p^* compared with the unidirectional flow test results. Increase solidity produces an increase in pressure drop.

Fig 6 shows a deviation of oscillating flow test results from the unidirectional flow test results when the power outputs are compared. Once again, within the scatter of the data the effect of increase in f on W^* is not clear but it is apparent that the effect of flow oscillations is to increase W^* at the lower solidity of 0.52 and decrease

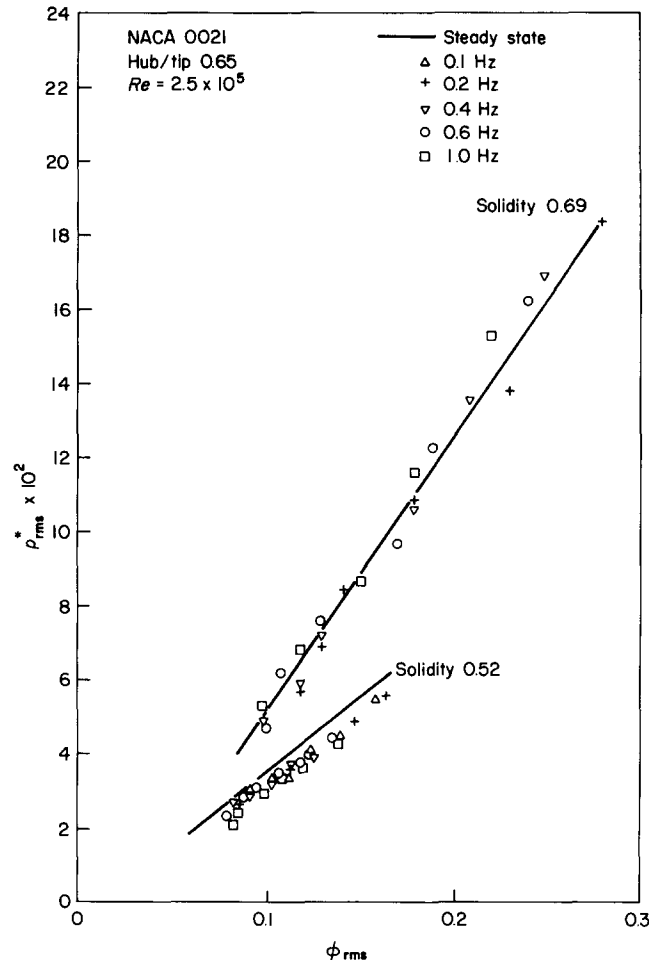


Fig 6 Effect of solidity on power output

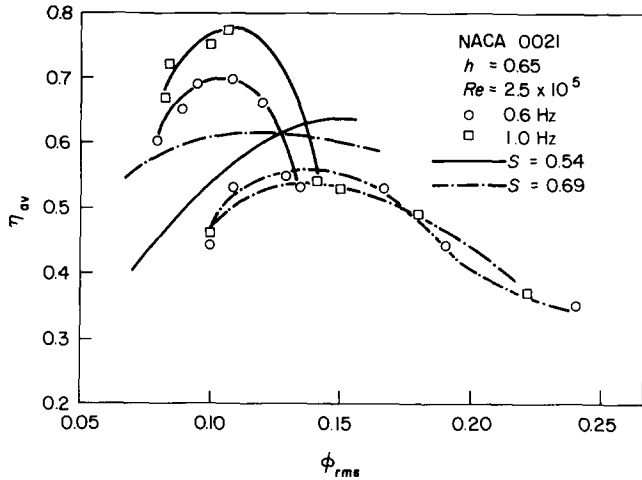


Fig 7 Effect of solidity on efficiency

W^* at a higher solidity of 0.69 when the results are compared with those at zero frequency. Stalling is very noticeable in this plot, with the increase in frequency promoting an earlier stall.

Increase in air flow frequency results in an improvement in efficiency at $S=0.52$ and reduction in efficiency at $S=0.69$ (Fig 7). The overall turbine solidity of 0.52 corresponds to a tip solidity of 0.43 where as $S=0.69$ corresponds to a tip solidity of 0.57. As a major part of energy conversion takes place nearer the tip rather than the hub, the tip solidity is a good representation of the turbine geometry. In this context, a tip solidity of 0.43 can be regarded as a low value and 0.57 as a high value. The results suggest that at low solidity the aerodynamic performance, which can be measured in terms of C_L/C_D , of the blades improves with the frequency of air flow oscillations. Although there are no wind tunnel tests available for a direct comparison, oscillating isolated aerofoil tests do show that the air flow oscillations produce hysteresis effects⁵ which tend to postpone separation and therefore increase C_L/C_D and hence increase the aerodynamic performance of an aerofoil blade.

Bringing the blades to a close proximity, however, seems to produce a deterioration in the aerodynamic performance of the blades in an oscillating air flow as the results of the high solidity tests tend to show. This is possibly due to the fact that the wake shed by a blade lowers the C_L and increases C_D on a blade just behind it⁶, ie on blades closely arranged in tandem cascade.

For a given frequency the efficiency of the turbine therefore reduces with the increase in solidity showing the reduction in aerodynamic performance of the turbine when the blades are brought closer to each other; efficiency increases when the blades are further apart.

The optimum pressure drop p_{opt}^* and the optimum power coefficient W_{opt}^* for maximum efficiency as a function of Strouhal number of air flow S_t , for both the solidities are shown in Fig 8. In this non dimensional plot it is interesting to note that for both solidities the pressure drop decreases with the Strouhal number. However, at the high solidity the effect of Strouhal number on p_{opt}^* appears to be small. The $W_{opt}^* - S_t$ curves show a continuous drop in W_{opt}^* with S_t for both the solidities. The results from the oscillating flow tests correlate very well with the results from the unidirectional flow tests.

Fig 9 shows the effect of Strouhal number on the maximum value of cyclic efficiency, η_{max} , and the corresponding flow coefficient ϕ_{opt} . Improvement in η_{max} obtained at higher S_t for the low solidity has already been explained. Typically at $S=0.52$ and $S_t=10^{-3}$ an increase in η_{max} of 5-6% is observed when compared to the value of η_{max} at $S_t=0$. There is a corresponding reduction in η_{max} of 5-6% at $S=0.69$. These results show that there is an advantage in using a turbine of low solidity in oscillating air flows. Such a turbine, however, can only operate at low pressure drops, as seen earlier, and may not be self starting². These curves also indicate that for $S=0.69$ the values of η_{max} may be independent of S_t for the operating region of the Wells turbine, that is $0 < S_t < 10^{-3}$. The optimum flow coefficient values ϕ_{opt} in Fig 9 show a larger sensitivity to S_t at the low values of S . This figure shows that the design value of ϕ must be somewhat lower than that based on unidirectional flow tests.

Starting torque

In order to understand the starting torque produced by an oscillating air flow, the turbine rotor was locked and the variation of torque τ at zero speed ω with pressure drop Δp and flow rate Q were measured at several air flow frequencies. Figs 10 and 11 show the variation of pressure drop with flow rate and torque with flow rate for two turbine solidities and frequencies of 0.2 Hz and 1 Hz and a constant average $R_x = 5.5 \times 10^4$.

The results could be expressed in non-dimensional

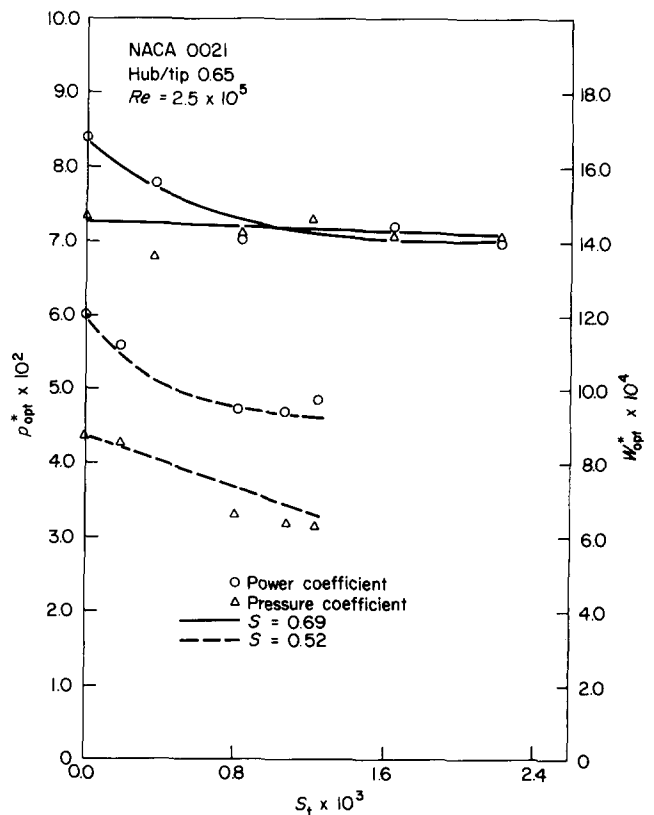


Fig 8 Effect of Strouhal number on optimum pressure drop and power output

$$\tau^* = \frac{\bar{\tau}}{\Delta p} \tag{7}$$

Strouhal number, S_x :

$$S_x = \frac{f_c}{V_x} \tag{8}$$

Reynolds number R_x :

$$R_x = \frac{U_c}{V_x} \tag{9}$$

It should be noted that the turbine blades are subjected to incidences ranging from $+90^\circ$ to -90° during these tests. Unpublished results from similar tests in a unidirectional air flow rig are also shown. The $\Delta p \sim Q$ curves show a quadratic relationship ($\Delta p = KQ^2$) for all frequencies including $f=0$. The rotor behaves like an orifice plate when it is stationary and the constant of proportionality K is a function of solidity and is virtually independent of frequency at these low frequencies.

Fig 11 shows that there is a negligible effect of low frequencies on torque produced at a given flow rate for $S=0.52$ and $S=0.69$. It has been suggested that the trends could change at low solidities and high frequencies.

The true effect of air flow frequency on the turbine blades can be observed from Fig 12 where a non-dimensional value of torque $\bar{\tau}$ is plotted against a non-dimensional value of pressure drop Δp for several values

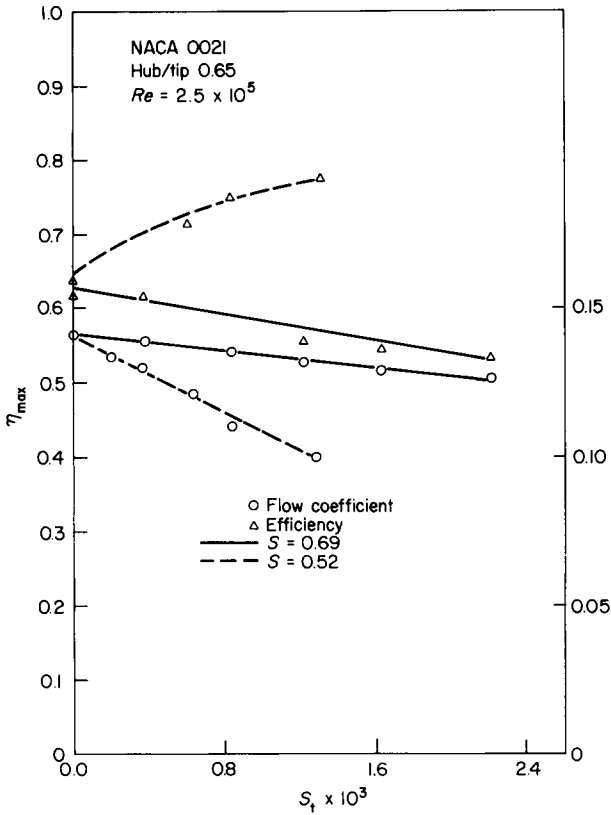


Fig 9 Effect of Strouhal number on efficiency and flow coefficient

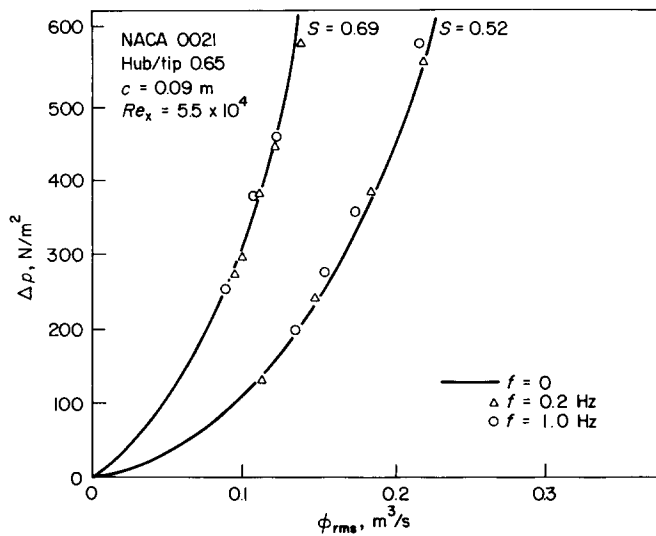


Fig 10 Variation of pressure drop with flow rate

form, viz pressure drop coefficient $\bar{\Delta p}$:

$$\bar{\Delta p} = \frac{\Delta p}{\frac{1}{2}\rho V_x^2}$$

reduced torque, $\bar{\tau}$:

$$\bar{\tau} = \frac{\tau}{\frac{1}{2}\rho V_x^2 n c D_{v/4} (1-h^2)}$$

torque coefficient, τ^* :

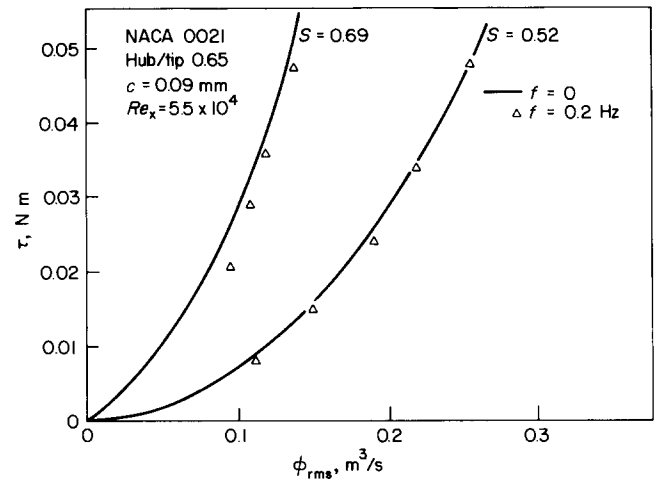


Fig 11 Variation of torque with flow rate

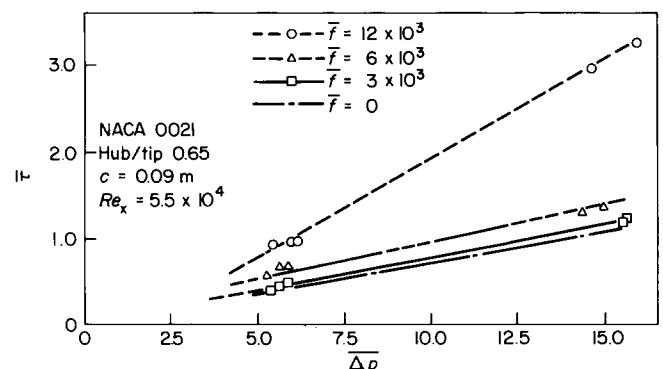


Fig 12 Variation of torque with pressure drop

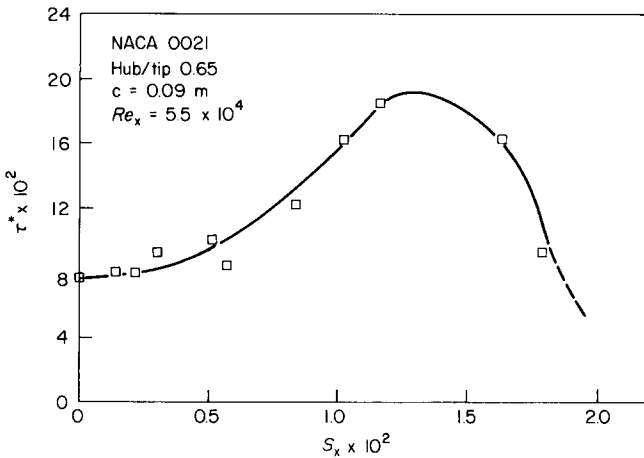


Fig 13 Variation of torque with Strouhal number

of Strouhal number S_x based on axial flow velocity V_x . For a given value of S_x , $\bar{\tau}$ varies linearly with Δp and this is independent of S . For a given Δp , $\bar{\tau}$ increases significantly with S_x . The values of $\bar{\tau}$ at $S_x = 12 \times 10^{-3}$ are twice those at $S_x = 0$. At starting, $\alpha = 90^\circ$ and hence the value of tangential force coefficient C_t and axial force coefficient C_x are fixed for a given S , R_x and profile. $\bar{\tau}$ is equal to C_t and Δp is a function of both h and S . Since the value of h is constant, the increase in $\bar{\tau}$ with Δp is associated with the increase in S .

A unique relationship between the torque and frequency can be obtained if a coefficient of torque τ^* is plotted against S_x as shown in Fig 13. This correlation is independent of solidity, pressure drop and flow rate and is a function of S_x only. A rapid increase in τ^* with S_x is noticeable from this figure but only up to a value of $S_x = 1.25 \times 10^{-2}$ beyond which there is a sharp reduction in τ^* . As most of the wave energy devices are likely to operate with $S_x < 10^{-2}$, the drop in τ^* with S_x at high

frequencies is not important. Therefore this curve suggests that considerably higher starting torque is achieved in oscillating air flow than that predicted from unidirectional flow tests.

Conclusions

The performance of the Wells self-rectifying air turbine in oscillating air flow is presented in this paper. The running performance in oscillating air flow is solidity dependent. Increase in solidity produces an increase in pressure drop and reduction in maximum efficiency. Increasing the frequency of air oscillation results in an increase in power output and efficiency for low solidity and vice versa for high solidity. The oscillating air flow produces a higher starting torque than a unidirectional air flow.

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